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BASELINE CHEMISTRY MODEL FOR PLUME RADIATION PREDICTIONS

Edward R. Fisher, et al

Physical Dynamics, Incorporated

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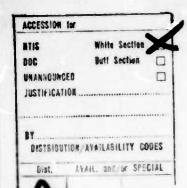
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INTRODUCTION

As part of the DARPA Plume Physics Program, extensive modeling calculations have been performed on both field and test environments. These modeling calculations generally involve fluid dynamical, chemical and radiative inputs, each of which may be subject to certain approximations in order to yield a final value for comparison with observation. The purpose of this paper is to outline four chemistry models which can be used for baseline predictions on the IR radiative characteristics of rocket exhaust plumes.

Although these models were initially generated for use in high altitude plume environments where continuum fluid approximations are still valid, the resulting chemistry is also useful to both lower altitude regimes and test environments. The four models developed to date consist of the following:

- (1) APOLLO model
- (2) RELAXATION model
- (3) HYDROCARBON model
- (4) AMINE model .

Each model consists of a chemical mechanism consisting of a sequence of chemical reactions and energy transfer processes with temperature dependent rate coefficients as well as the radiative lifetimes of appropriate species in the IR region of the spectrum. A preliminary attempt to outline a UV model

has been documented elsewhere [Kummler, Fisher, Boynton (1972)].

The rationale behind the development of these baseline chemistry models stems from the desire to provide a common chemistry for direct comparison between different flow field models. In addition, these baseline models provide a state-of-the-art chemistry package to use in making reasonable predictions of plume radiation as well as the radiation observed from rocket test plumes. As new chemical reactions and improved rate coefficient information becomes available, the chemistry models will be updated.

As will become clear in reviewing the chemistry models, many rate coefficients are poorly known. However, the attempt in these models has been to include as much of the important chemistry as is feasible with rate coefficients estimated either from theory or by comparison with other better known rate processes. The output from model calculations can then provide the basis for the design of experiments aimed at quantifying the important chemical rates. This has already been done in several cases, as will be discussed later.

The chemical processes included in these baseline chemistry models must reflect the highly non-equilibrium state of high altitude plumes. Therefore, the following types of chemistry processes have been included in the current models:

Chemi-excitation processes - these are chemical reactions in which the exothermicity of the reaction is either partly or totally deposited directly into an internal energy state of a product molecule. Generally speaking, interest has been focused on product vibrational states since these excited species can directly or indirectly (by transferring the vibrational energy to another species or mode) contribute to observed IR radiation levels. An example of this type of process is

он +
$$H_2 \longrightarrow H_2 O^{\ddagger} + H$$

where ${\rm H_2O}^{\ddagger}$ represents a vibrationally excited molecule. Sufficient reaction energy exists in this process to excite either the stretching or bending modes. As a result of interest generated in this and related processes, experiments have been performed on the quantum yield of these reactions [Trainor (1973)].

Relaxation processes - these are energy transfer processes in which vibrational energy is exchanged with translational or rotational energy (VT) in a collision, or vibrational energy is exchanged directly between vibrational modes (VV). As an example of the former case (VT), the process

$$co_2^{\dagger} + o \longrightarrow co_2 + o$$

is thought to be important in the mixing layer of high altitude plumes. Since vibrational relaxation processes involving atomic species are known to be important in cases where direct experiment has been performed, plume radiation by 0 atoms is believed to be of major importance in the mixing layer of high altitude plumes. The rates for several of these 0 atom relaxation processes have

been measured as part of the Plume Physics Program and the results are included in the models. As an example of a VV process

$$N_2^{\dagger} + H_2^{\circ} \longrightarrow N_2 + H_2^{\circ}(v_2)$$

may contribute to the excitation of H_2O by atmospheric N_2 which is vibrational excited in passing through the air shock prior to entering the mixing layer. Non resonance in vibrational energy in VV processes is supplied or absorbed by the local kinetic and rotational energy depending on the direction of the reaction. For both VT and VV processes, a significant amount of rate coefficient information was obtained from the results of past laser studies.

Excited state processes - these are chemical reactions in which
the reactants are in excited internal energy states. By
detailed balance arguments based on chemi-excitation reactions, excited state processes are expected to be important in the plume environment. A particularly significant excited state process is

$$O + H_2^{\dagger} \longrightarrow OH + H$$
.

This reaction between atmospheric O atoms and vibrationally excited $\rm H_2$ may be the source reaction for the vibrationally excited OH observed in APOLLO at high altitudes. This reaction also generates ground state OH which can subsequently react with more $\rm H_2$ to generate $\rm H_2O$ as outlined under chemi-excitation reactions above. The ground state reaction

$$O + H_2 \longrightarrow OH + H$$

has been extensively measured and is well known over the temperature range of interest in plume modeling. The activation energy of 9.5 kcal/mole represents a substantial barrier to the generation of OH in low temperature environments such as rocket plumes. However, if the H₂(v=1) state, with an energy of 13 kcal/mole above H₂(v=0), can provide this activation energy then the excited state reaction might proceed with near gas kinetic rate providing an important production reaction for OH. Unfortunately, this excited state reaction is not easily amenable to experimental determination, but current theoretical efforts underway at Wayne State University in collaboration with Brookhaven National Laboratory may permit evaluation of the rate and energy flow in this excited state reaction.

thermic chemical reactions between atmospheric O atoms and partially reacted or unreacted fuel fragments which find their way into the plume gas through boundary layer flow, peripheral flows (turbine and exhausterator flows), or incomplete combustion processes. Although the fraction of fuel fragments in the plume flow is generally small (of order 1%) the rate and energetics of fuel species make even this small amount significant in an overall radiation energy balance. An example of a fuel afterburning reaction is the the reaction

$$O + N_2H_4 \longrightarrow H_2O + N_2H_2$$

which proceeds very rapidly with an exothermicity of about 90 kcal/mole. As discussed in a later section of

this report, the H₂O molecule is likely to be highly vibrationally excited and this provides a substantial source of IR radiation.

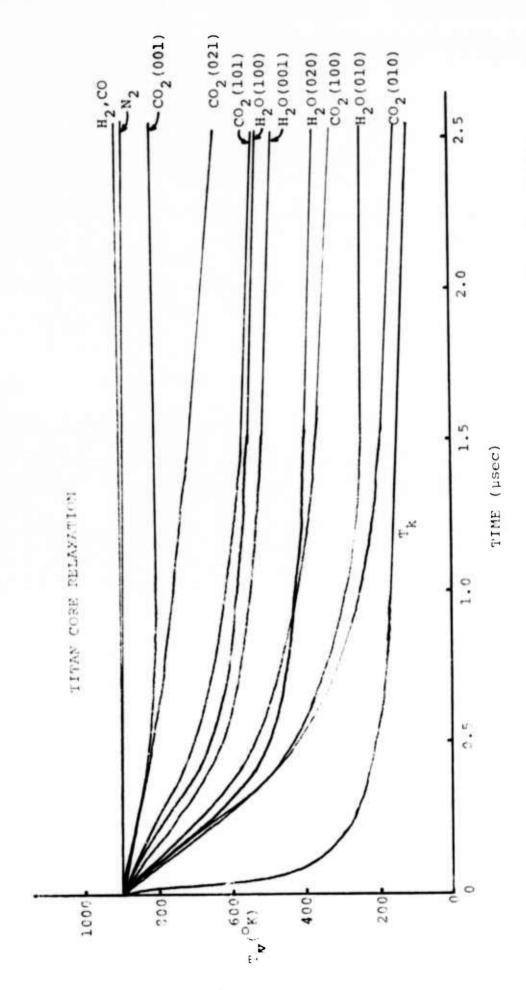
DISCUSSION OF MODELS

It is not the intention of this paper to provide an exhaustive review of the four plume chemistry models, but to indicate their availability, to provide some measure of their reliability, and to present some preliminary calculations performed using these models. Although reports have been prepared detailing the four models [Fisher (1973a); Kummler, Fisher, Boynton (1972, 1973)], a summary tabulating the reactions, rate coefficients and references are given in Appendix I.

Relaxation Model

In this section, we will present a calculation using the RELAXATION model on a TITAN core streamline. This calculation will indicate the detail included in the model as well as the methodology of reducing the size of the chemistry package for specific applications. A report detailing these TITAN core calculations is being prepared [Fisher (1973b)].

In Figure 1, the effective vibrational temperatures of important modes on an outermost TITAN core streamline are shown.



core flow calculated using the RELAXATION Model [Fisher (1973b)]. Vibrational temperatures in an "outermost" streamline of a TITAN FIGURE 1.

The calculation was done using a Runge-Kutta-Merson integration procedure with algebraic subroutine for treating the chemistry which is in a quasi steady state. These quasi steady state reactions generally force the system of chemical reactions into a "stiff" condition which greatly increases integration time for standard integrations packages. The algebraic subroutine provides a mechanism to bypass this complication as has been extensively discussed before [Fisher, Kummler (1968)].

The general features that emerge from this calculation are that the diatomic molecules are vibrationally frozen at rear exit conditions; the bending modes of ${\rm CO_2}$ and ${\rm H_2O}$ are nearly in equilibrium with the translational temperature while the stretching modes are partially frozen depending on the extent of VV coupling to the diatomic species. Due to the rapid VV transfer with ${\rm N_2(v=1)}$, ${\rm CO_2(001)}$ is rather vibrationally hot along this streamline. The modes of ${\rm CO_2}$ which depend on ${\rm CO_2(001)}$ as a source, i.e., ${\rm CO_2(021)}$ and ${\rm CO_2(101)}$ are also rather vibrationally hot but reflect the fact that the bending modes of ${\rm CO_2}$ are rather vibrationally cold.

The RELAXATION model maintains the Fermi resonance between ${\rm CO_2(100)}$ and ${\rm CO_2(020)}$ as a VT process with nearly gas kinetic efficiency independent of the collision species, while the ${\rm H_2O(100)}$ and (001) modes are assumed to be in quasi equi-

librium at the local kinetic temperature. This latter assumption stems from the closeness in energy resonance between these modes as well as the sparcity of data regarding the relaxation of higher vibrational levels in ${\rm H}_2{\rm O}$.

The code used in these calculations also provides a "picture" of the important rate processes involving each excited species. This output, called HISTORY, is provided at arbitrary time intervals in the calculation and permits evaluation of the significant reactions both supplying and destroying any particular level. Two such reaction profiles are shown in Tables I and II for two times. Note that the rate coefficients are not generally the same for the two tables reflecting the change in temperature while the difference in rates reflect not only the change in excited species density due to relaxation but also due to pressure. Both the translational temperature and pressure are printed at the top of the respective tables. By using these rate profiles, it is anticipated that the number of relaxation processes can be reduced for subsequent calculations in similar environments.

Since the hydrocarbon afterburning model has been discussed previously [Kummler, Fisher, Boynton (1972)], the amine afterburning model will be more extensively detailed here. Further background information can be obtained in report form [Kummler, Fisher, Boynton (1973)].

TABLE I

Chemical reactions, rate coefficients and separate rates at beginning of relaxation calculation in TITAN core flow as obtained from subroutine HISTORY [Fisher (1973b)].

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51.	* 00	+	0	11	00	+	0	.759E-1	•	
52.	00	+	0	H	*00	+	0	.512E-1	•	
53.	* 00	+	Н2	11	00	+	Н2	.947E-1	.841	-
54.	00	+	H2	11	*00	+	Н2	.254E-1	.795	
55.	*00	•	H20	11	00	+	H20	3.947E-14	5.636E	19
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72.	0	-	Н2	II	C	C	H2	.0C2E-1	.612	
73.		C.	H20	H	0	0	N	.947E-1	•	19
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TABLE II

Chemical reactions, rate coefficients and individual rates at 1.62 μ sec in TITAN core flow as obtained from subroutine HISTORY [Fisher (1973b)].

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43.	H20	0	0	H	N	20	0	PE-1		
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79. CO2 001+ H2	1. CO2 561+ H2	2. CO2 "3"+ H2	3. CO2 (41+ C	5. 002 011+ 00	6. CO2 37+ C	7. CO2 001+ N	8. CO2 C3C+ N	9. H2* + H20 SC	A. H2 + H2	1. H2* + OH	2. H2 + OH*	3. N2* + H20 91	4. N2 + H20 CO	5. N2* + H20 n	6. N2 + H20 01	7. H20 C10+ H20 61	8. HZO 320+ HZO 3C	9. CO2 101+ H20 00	30. CO2 010+ H20 01	€1. N2* + C02 n	r2. N2 + C02 (73. N2* + CO	04. N2 + C	05. CO2 001+ CO	6. CO2 OCT+ CO	e7. CO2 101+ H20 69	08. CO2 900+ H20 0€	09. CO2 C21+ H20 AC	19. CO2 nor+ H20 CC	11. CO2 101+ CO2 NO	12. CO2 10C+ CO2 07	13. CO2 020+ C	4. C02 010+ C02 01	5. H20 G2C+

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Chemical Models: Substituted Hydrazine

The chemistry of RHNNR'R" compounds is quite different from the hydrogen and hydrocarbon chemistry reviewed previously [Kummler, Fisher, Boynton (1972)] because of the characteristic strong NN bonding. In addition, the different chemical behavior of unsymmetrical dimethyl hydrazine (UDMH),

$$H_2NN(CH_3)_2$$
,

and hydrazine or monomethyl hydrazine (MMH),

precludes some of the simplicity of the generalizations inherent in the lumped parameter approach to the hydrocarbon oxidation.

Major radiating steps in the plume appear to be associated with the initial attack of 0 atoms on the fuel-like species, rather than the secondary reactions of products of the initial attack as in hydrocarbon oxidation.

Oxygen atom attack on hydrazine has been studied using high intensity crossed molecular beams [Gehring, et al (1970), Foner and Hudson (1970)]. The products of the reaction have been found to be N_2H_2 + H_2O , via the elementary reaction:

$$O + H N - N H \rightarrow HN = HN + H_2O.$$

This reaction is highly exothermic (90 kcal) and forms a completely new water molecule by simultaneously abstracting two hydrogen atoms. The expected (from, for example, the analogous O attack on ethylene) atom transfer reaction,

$$O + \frac{H}{H}N - N \frac{H}{H} \rightarrow N_2 H_3 + OH ,$$

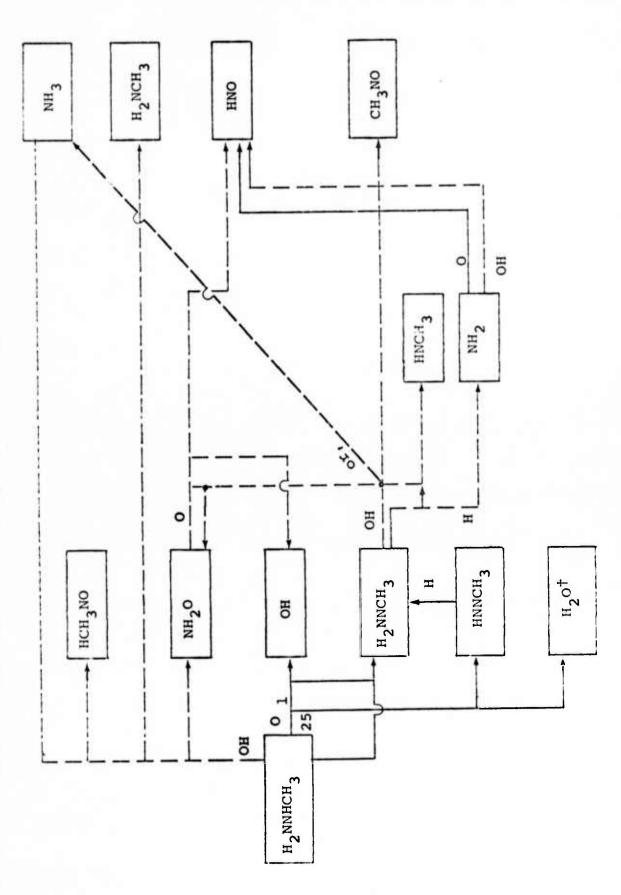
which is 26 kcal exothermic, is about 1/25 as probable [Foner and Hudson (1970)]. The reaction of O + UDMH can only produce OH.

Once significant quantities of OH become available through the latter reactions, additional processes are possible. Again, high intensity molecular beams [Gehring, et al (1970)] have been used to identify the products of the reaction:

$$OH + N2H4 \rightarrow NH3 + NH2O ,$$

which is the key reaction in severing the N-N bond.

The likely mechanism for oxidation of substituted hydrazine compounds is illustrated schematically in Figure 2. The reaction rate constants for the principal reactions are given in Table III. A more complete literature review of pertinent reactions has been given by Kummler, Fisher and Boynton (1973), while a summary of the Amine model is given in Appendix I.



Simplified O atom oxidation mechanism for MMH (UDMH and N₂H₄ are similar) illustrating the stable species which can form in this system as intermediates. The dotted lines represent reactions which are probable, but which have not yet been experimentally observed. FIGURE 2.

TABLE III
INFRARED HYDRAZINE MODEL A

FUEL CHEMI-EXCITATION REACTIONS

Reaction	Assumed Products*	Forward Rate Constant cc/sec	Reaction Energy kcal/gmole	Reference
O + MMH → NH ₂ CH ₃	+ H ₂ O(001) (100) H ₂ O(020) (010)	2 × 10 ⁻¹¹	-102	Foner and Hudson (1970)
$O + N_2H_4 \rightarrow N_2H_2 +$	H ₂ O(001) (100) H ₂ O(020) (010)	1.8 × 10 ⁻¹¹	-90	Gehring, <u>et al</u> (1969)
$O + UDMH \rightarrow HN_2(CH_3)_2 + OH$		1×10^{-12}	-25	Analogy
$OH + N_2H_4 \rightarrow NH_3 + NH_2O$		2 × 10 ⁻¹¹	?	Gehring, <u>et al</u> (1970)

^{*} This puts about 40% of the reaction energy into the stretching modes of product H₂O. (For purposes of computational convenience, we have treated a molecule excited to the second level as two singly-excited molecules; the quenching is thus somewhat more efficient than it should be.)

The salient features of Figure 2 have been examined experimentally in the Wayne State University discharge flow sys-Comparitive results for $O + N_2H_4$ chemiluminescence are presented in Figure 3. The reaction of $\mathrm{H} + \mathrm{NO}_2$ produces OH , but is not sufficiently exothermic to produce more than v=2 in the product molecule. Hence, the OH fundamental between 2.7µ and 3.0 μ is broader than the overtone at 1.5 μ . (The instrumental response has not been taken out of Figure 2, so that the overtone appears larger than the fundamental. Calibration with the NO +O glow reverses this effect and gives excellent agreement with observed values of the overtone to fundamental intensity ratio.) For the reaction of O + C_2H_4 the overtone shows hot bands consistent with the greater available energy. The CO overtone at 2.3µ also appears and the intensity ratio of the overtone to fundamental ratio under these conditions indicates that some ${\rm H}_2{\rm O}(100)$ and ${\rm H}_2{\rm O}(001)$ radiation may be present. The reaction of $O + N_2H_4$ exhibits a much broader peak at 2.7 μ indicative of ${\rm H_2O}(100)$ and (001) radiation and a band at 1.9μ also indicative of ${\rm H_2O}$ emission is present. The emission of these systems will be analyzed in detail elsewhere.

In order to assess the potential magnitude of the resulting infrared emission, we have taken the mechanism of Figure 2 and reduced it to the simplest form consistent with the production and loss of vibrationally excited H₂O. We have also

included a number of VT and VV excitation and de-excitation reactions for H₂O in order to have a basis for comparing calculations with and without chemi-excitation and to provide for quenching of excited species produced by reaction. The mechanism is presented in Appendix I; rate data for the H₂O VV and VT processes are the same as we have used in calculations on the Apollo system [Boynton (1973)] and follow the recommendations of the Plume Chemistry Workshop held at Physical Dynamics, Berkeley, in January 1973 [Fisher (1973a)].

The rate equations for the above mechanism were integrated through a flow field simulating the mixing layer of a Titan II exhaust plume at 120 km. Details of the flow field calculation will be presented in a separate report, but a brief outline is presented here. The exit plane conditions are taken from an unpublished nozzle flow calculation by D. Thomas (Aerospace Corp.) who included finite-rate chemistry and film cooling effects. External flow calculations were performed with the MULTITUBE and BOW codes [Boynton (1971)]. Unburned fuel concentrations of 5% in the film cooling layer and 1% in the interior exhaust gases were assumed; reactions of N₂H₄ and UDMH with 0 were included in the flow field calculation, so that depletion of these species would be properly accounted for in a subsequent (overlay) chemistry calculation.

Figure 3. A comparison of the infrared emission spectra for the reactions $0+N_2H_4$, $0+C_2H_4$ and $H+NO_2$. Concentrations of 0 and H were approximately 10μ and the concentrations of N_2H_4 , C_2H_4 and NO_2 were approximately 2μ , 10μ and 4.5μ respectively. All spectra were taken at a contact time of less than 2 msec. Both 0 and H were diluted in a stream consisting primarily of He.

- H + NO $_2$ \rightarrow ; 2000μ slits, 50μ vFS, 100μ FS, τ = 10, 1 RPM, $\frac{5 \text{ min}}{\text{in}}$, 800μ He, 200μ H $_2$, 5μ NO $_2$, 5μ H, 1st window.
- $O + C_2H_4 \rightarrow$; 1000 μ slits, 50 μ vFS, 100 μ FS, τ = 10, 1 RPM, $\frac{5 \text{ min}}{\text{in}}$, 10μ C₂H₄, 800 μ He, 200 μ O₂, 1st window.
- $O + N_2H_4 \rightarrow$; 2000 μ slits, 20 μ vFS, 100 μ FS, $\tau = 3$, 4 RPM, $\frac{2 \text{ min}}{\text{in}}$ $2\mu \ N_2H_4$, 900 μ He, 90 μ O₂, 11 μ O, 1st window.

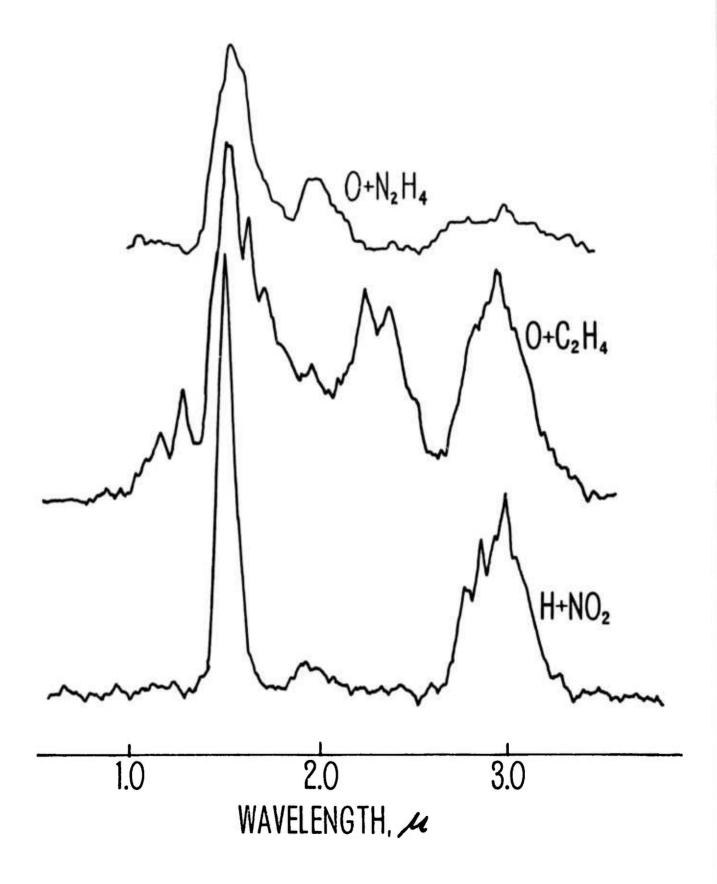


FIGURE 3.

Results of these exploratory calculations are shown in Figure 4. We show cumulative radiant intensity as a function of distance behind the vehicle for the sum of the $\rm H_2O$ (100) and (001) bands and the OH band (which is negligible in this calculation because important formation mechanisms such as $\rm H_2 + O$ \rightarrow OH + H have not been included). The lower curve is for the $\rm H_2O$ excitation processes alone, and the upper curve includes the effects of the simplified chemi-excitation mechanism.

With the assumption of product energy distribution together with those about initial fuel concentrations, there is a substantial difference in calculated radiant intensity with and without chemiluminescent reaction.

These calculations may overestimate the contribution of chemiluminescence to the plume signature for any or all of the following reasons:

1) The exhaust-gas concentrations of N_2H_4 and UDMH may be too high. In that connection, note that N_2H_4 (the major reacting species in this system) is a good mono-propellant and decomposes spontaneously at elevated temepratures to give less reactive NH_3 , N_2 , and H_2 . The amount of external reaction may depend on the extent of such decomposition as well as the overall unburned fuel level. This should not be the case with MMH, which is more stable.

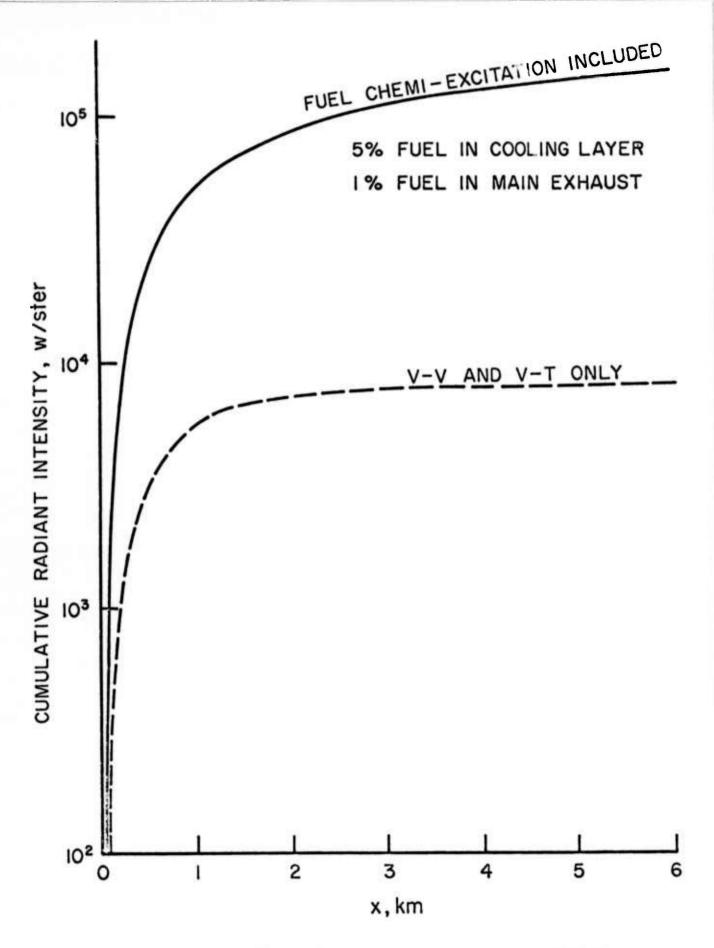


FIGURE 4. The effect of chemi-excitation on predicted cumulative radiant intensity.

- The photon yields may be optimistic, although a rather drastic decrease would still allow substantial chemi-excitation relative to VT and VV excitation.
- 3) Unburned fuel could be ejected from the engine or turbine exhaust in the form of droplets which would have to be vaporized in the plume before reacting.

 [Webber, et al (1972), using a detailed combustion code to evaluate the external surface contamination potential due to unburned propellants from two small engines, indicates that essentially all unburned propellant leaving these engines is in the form of small droplets.]

However, even with these qualifications, it is evident that there is a distinct possibility that external chemiluminescence of unburned fuel species comprises a substantial fraction of aminefueled missile plume signatures. This possibility should not be ignored in studies aimed at elucidating plume emission mechanisms.

SUMMARY

The availability of baseline chemistry models for four rocket plume environments has been established. The mechanism for determining the important relaxation processes in systems involving non-equilibrium distributions of N_2 , H_2 , CO, O, H, H_2O and CO_2 has been presented and preliminary calculations shown. An amine/substituted hydrazine model has been detailed and calculations appropriate to the TITAN mixing layer presented. As was found in the case of hydrocarbon afterburning in the ATLAS vehicle [Fisher (1973c)], minor quantities of unreacted amine or hydrazine in the plume flow can lead to a significant amount of plume radiation in the IR.

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APPENDIX I

BASELINE CHEMISTRY MODELS

(1)	APOLLO Model	34
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(1) APOLLO MODEL

	Reaction	Reference	Rate Constant (cc/molecule-sec)
1)	O + H ₂ → OH + H	(a,b)	$1.65^{-13} \exp(-2500/T)$
2)	H + OH → O + H ₂	(c)	$6.9^{-14} \exp(-1425/T)$
3)	$O + H_2(v=1) \rightarrow OH + H$	(d)	$1.3^{-10} \alpha (\alpha = 0.9, 0.5)$
4)	$H + OH \rightarrow O + H_2(v=1)$	(c)	$5.5^{-11} \alpha \exp(-4925/T)$
5)	$O + H_2(v=1) \rightarrow OH(v=1) + H$	(d)	$1.3^{-10} (1-\alpha)$
6)	$H + OH(v=1) \rightarrow O + H_2(v=1)$	(c)	5.5^{-11} $(1-\alpha) \exp(215/T)$
7)	$OH + H_2 \rightarrow H_2O + H$	(d)	$3.8^{-11} \exp(-2600/T)$
8)	$H + H_2O \rightarrow OH + H_2$	(c)	$1.5^{-10} \exp(-10075/T)$
9)	$OH + H_2(v=1) \rightarrow H_2O + H$	(e)	1.5-10
10)	$H + H_2O \rightarrow OH + H_2(v=1)$	(c)	$5.8^{-10} \exp(-13475/T)$
11)	$OH(v=1) + H_2 \rightarrow H_2O + H$	(e)	1.5 ⁻¹¹
12)	$H + H_2O \rightarrow OH(v=1) + H_2$	(c)	$5.8^{-11} \exp(-12475/T)$
13)	OH + OH → H ₂ O + O	(a)	$9.6^{-12} \exp(-390/T)$
14)	O + H ₂ O → OH + OH	(a)	$9.5^{-11} \exp(-9000/T)$
15)	O + OH → O ₂ + H	(f)	4-11
16)	$H + O_2 \rightarrow O + OH$	(c)	$6.9^{-10} \exp(-8400/T)$
17)	$O + OH(v=1) \rightarrow O_2 + H$	(f)	4-11
18)	$H + O_2 \rightarrow O + OH(v=1)$	(c)	$6.9^{-10} \exp(-13400/T)$

ENERGY TRANSFER REACTIONS*

Vibration-Translation Processes:

19)
$$H_2(v=1) + H + H_2 + H$$
 (g) $3.2^{-11} \exp(-1400/T)$ 20) $H_2(v=1) + 0 + H_2 + 0$ (e) $3.0^{-13} \sqrt{T}$ 21) $H_2(v=1) + H_2 + H_2 + H_2$ (h,i) $1.5^{-7} \exp(-139/T^{1/3})$ 22) $H_2(v=1) + H_20 + H_2 + H_20$ (e,j) $1.5^{-7} \exp(-139/T^{1/3})$ 23) $N_2(v=1) + H + N_2 + H$ (u) $1.9^{-6} \exp(-164/T^{1/3})$ 24) $N_2(v=1) + 0 + N_2 + 0$ (k,1,m) $1.2^{-13} \exp(-23/T^{1/3})$ 25) $N_2(v=1) + H_2 + N_2 + H_2$ (n) $3.9^{-8} \exp(-164/T^{1/3})$ 26) $N_2(v=1) + H_20 + N_2 + H_20$ (e,j) $3.9^{-8} \exp(-164/T^{1/3})$ 27) $OH(v=1) + H + OH + H$ (e) 3^{-11} 28) $OH(v=1) + O + OH + O$ (e) 3^{-11} 29) $OH(v=1) + H_2 + OH + H_2$ (e) 1^{-11} 30) $OH(v=1) + H_20 + OH + H_20$ (e) 3^{-12} 31) $H_2O(010) + H + H_2O(000) + H$ (e) 1^{-11} 32) $H_2O(010) + O + H_2O(000) + O$ (e) 1^{-11} 33) $H_2O(010) + H_20 + H_2O(000) + H_2$ (e,o) 1^{-11} 35) $H_2O(020) + H_20 + H_2O(010) + H$ (t) 2^{-11} 36) $H_2O(020) + O + H_2O(010) + O$ (t) 2^{-11} 37) $H_2O(020) + H_2 + H_2O(010) + H_2$ (t) 2^{-11} 38) $H_2O(020) + H_20 + H_2O(010) + H_2$ (t) 2^{-11} 38) $H_2O(020) + H_20 + H_2O(010) + H_20$ (t) 2^{-11} 38) $H_2O(020) + H_20 + H_2O(010) + H_20$ (t) 2^{-11}

^{*} Note that the reverse rate coefficient can be calculated directly from detailed balance.

39)
$$H_2O(001) + H \rightarrow H_2O(000) + H$$
 (e)

40)
$$H_2O(001) + O \rightarrow H_2O(000) + O$$
 (e) 1^{-11}

41)
$$H_2O(001) + H_2 \rightarrow H_2O(000) + H_2$$
 (o) 2^{-13}

42)
$$H_2O(001) + H_2O \rightarrow H_2O(000) + H_2O$$
 (q) 1^{-12}

43)
$$H_2O(001) + H \rightarrow H_2O(020) + H$$
 (e) 1^{-11}

44)
$$H_2O(001) + O \rightarrow H_2O(020) + O$$
 (e) 1^{-11}

45)
$$H_2O(001) + H_2 + H_2O(020) + H_2$$
 (e) 1^{-11}

46)
$$H_2O(001) + H_2O \rightarrow H_2O(020) + H_2O$$
 (e) 1^{-11}

Vibration-Vibration Processes

47)
$$H_2(v=1) + H_2O(000)$$

 $+ H_2(v=0) + H_2O(001)$ (r) $2.5^{-14} T^{5/6} \exp(-30.1/T^{1/3})$

1-11

48)
$$H_2(v=1) + OH(v=0)$$

 $\rightarrow H_2(v=0) + OH(v=1)$ (e) 1^{-12}

49)
$$N_2(v=1) + H_2O(010)$$

 $\rightarrow N_2(v=0) + H_2O(001)$ (e) 3^{-13}

50)
$$N_2(v=1) + H_2O(000)$$

 $+ N_2(v=0) + H_2O(010)$ (s) 9.4⁻¹¹ $exp(-64/T^{1/3})$

51)
$$H_2O(010) + H_2O(010)$$

 $+ H_2O(020) + H_2O(000)$ (e) $1^{-12} \sqrt{T}$

52) $\Pi_2^{\rm O}(001)$ and $\Pi_2^{\rm O}(100)$ assumed in quasi-equilibrium at the local translational temperature.

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- e) estimate
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- h) J.H. Kieffer and R.W. Lutz, J. Chem. Phys. <u>44</u>, 658(1966) <u>45</u>, 3888(1966).
- i) C. Joffrin, J. Ducuing and J.P. Coffinet, Opt. Commun. 2, 245(1970).
- j) past measurements involving water presumably measure both VV and VT processes.
- k) W.D. Breshears and P.F. Bird, J. Chem. Phys. 48, 4768 (1968).
- R.J. McNeal, M.E. Whitson Jr. and G.R. Cook, preprint of work submitted to Chem. Phys. Letters (1972). Aerospace Report SPL 3303, July (1972).
- m) D.J. Eckstrom, Stanford Research Institute, unpublished results.
- n) D.R. White, J. Chem. Phys. 46, 2016(1967).
- o) estimate based on Marriott's cross section values see E.R. Fisher, AFCRL Special Report, 72-0539, August(1972).
- p) estimate based on all available data
- q) estimate based on analogy to HF VT relaxation

- r) estimate based on the modified Rapp model, J. Chem. Phys. 40, 573(1964); Fisher unpublished results.
- s) estimate based on data of C.W. Von Rosenberg, K.N.C. Bray and N.H. Pratt, J. Chem. Phys. <u>56</u>, 3230(1972) and 13th Symp. (Intern.) on combustion p. 89,(1971).
- t) scaled from the analogous rate constant for exciting the lowest bending mode.
- u) estimated based on the rate for CO + H; see reference (o) under RELAXATION MODEL (2).

(2) RELAXATION MODEL

Vibr	ation-Translation Processes	Reference	Rate Constant (cc/molecule-sec)
1)	$H_2(v=1) + H + H_2 + H$	(a)	$3.2^{-11} \exp(-1400/T)$
2)	$H_2(v=1) + O \rightarrow H_2 + O$	(b)	$3.0^{-13} \sqrt{T}$
3)	$H_2(v=1) + H_2 \rightarrow H_2 + H_2$	(c,d)	$1.5^{-7} \exp(-139/T^{1/3})$
4)	$H_2(v=1) + H_2O + H_2 + H_2O$	(b,e)	$1.5^{-7} \exp(-139/T^{1/3})$
5)	$N_2(v=1) + H \rightarrow N_2 + H$	(y)	$1.9^{-6} \exp(-164/T^{1/3})$
6)	$N_2(v=1) + 0 \rightarrow N_2 + 0$	(f,g,h)	$1.2^{-13} \exp(-23/T^{1/3})$
7)	$N_2(v=1) + H_2 \rightarrow N_2 + H_2$	(i)	$3.9^{-8} \exp(-164/T^{1/3})$
8)	$N_2(v=1) + H_2O + N_2 + H_2O$	(b,e)	$3.9^{-8} \exp(-164/T^{1/3})$
9)	$OH(v=1) + H \rightarrow OH + H$	(b)	3-11
10)	$OH(v=1) + O \rightarrow OH + O$	(b)	3 ⁻¹¹
11)	$OH(v=1) + H_2 \rightarrow OH + H_2$	(b)	1-11
12)	$OH(v=1) + H_2O \rightarrow OH + H_2O$	(b)	3 ⁻¹²
13)	$H_2O(010) + H \rightarrow H_2O(000) + H$	(b)	1-11
14)	$H_2O(010) + O \rightarrow H_2O(000) + O$	(b)	1-11
15)	$H_2O(010) + H_2 \rightarrow H_2O(000) + H_2$	(b,j)	1-11
16)	$H_2O(010) + H_2O + H_2O(000) + H$	₂ 0 (k)	1-11
17)	$H_2O(001) + H + H_2O(000) + H$	(b)	1-11
18)	$H_2O(001) + O \rightarrow H_2O(000) + O$	(b)	1-11
19)	$H_2O(001) + H_2 \rightarrow H_2O(000) + H_2$	(j)	2 ⁻¹³
20)	$H_2O(001) + H_2O \rightarrow H_2O(000) + H$	20 (1)	1-12

21)
$$H_2O(001) + H + H_2O(020) + H$$
 (b) 1^{-11}

22) $H_2O(001) + O + H_2O(020) + O$ (b) 1^{-11}

23) $H_2O(001) + H_2 + H_2O(020) + H_2$ (b) 1^{-11}

24) $H_3O(001) + H_2O + H_2O(020) + H_2O$ (b) 1^{-11}

25) $CO(v=1) + H + CO(v=0) + H$ (o) $5.3^{-7} \exp(-119/T^{1/3})$

26) $CO(v=1) + O + CO(v=0) + O$ (p) $1^{-8} \exp(-96/T^{1/3})$

27) $CO(v=1) + H_2 + CO(v=0) + H_2$ (q) $9^{-9} \exp(-119/T^{1/3})$

28) $CO(v=1) + H_2O + CO(v=0) + H_2O$ (b) $9^{-9} \exp(-119/T^{1/3})$

29) $CO_2(010) + H + CO_2(000) + H$ (b) $3.8^{-12} \exp(-7/T^{1/3})$

30) $CO_2(010) + O + CO_2(000) + O$ (p) $2.3^{-9} \exp(-77/T^{1/3})$

31) $CO_2(010) + H_2O + CO_2(000) + H_2O$ (s) 1^{-11}

33) $CO_2(010) + H_2O + CO_2(000) + H_2O$ (s) 1^{-11}

34) $CO_2(010) + CO_2 + CO_2(000) + CO_2O$ (r) $4.6^{-10} \exp(-77/T^{1/3})$

35) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $1^{-8} \exp(-208/T^{1/3})$

36) $CO_2(001) + O + CO_2(000) + O$ (t) $1^{-8} \exp(-208/T^{1/3})$

37) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

38) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + H_2O + CO_2(000) + H_2O$ (t) $9^{-9} \exp(-119/T^{1/3})$

39) $CO_2(001) + O + CO_2(000) + O$ (u) $4.6^{24} T^{-5.89} \exp(-4223/T -6.74 \exp(-3013/T -6.74 \exp(-$

41)
$$CO_2(001) + H_2O + CO_2(030) + H_2O$$

(v)
$$4.71^{-40}T^{4.54} \exp(2081/T + 454/T^{1/3} - 1729/T^{2/3})$$

42)
$$CO_2(001) + CO_2 + CO_2(030) + CO_2$$

(v)
$$9.16^{23}T^{-5.89} \exp(-4223/T - 672.7/T^{1/3} + 2683/T^{2/3})$$

43)
$$co_2(001) + co + co_2(030) + co$$

(v)
$$6.87^{23}T^{-5.59} \exp(-4223/T -672.7/1.1/3 + 2683/T^{2/3})$$

44)
$$CO_2(001) + N_2 + CO_2(030) + N_2$$

(v)
$$6.87^{23}T^{-5.89} \exp(-4223/T)$$

 $-672.7/T^{1/3} + 2683/T^{2/3}$

Vibration-Vibration Processes

45)
$$H_2(v=1) + H_2O(000)$$

 $+ H_2(v=0) + H_2O(001)$

(m)
$$2.5^{-14} T^{5/6} \exp(-30.1/T^{1/3})$$

46)
$$H_2(v=1) + OH(v=0)$$

 $\rightarrow H_2(v=0) + OH(v=1)$

47)
$$N_2(v=1) + H_2O(010)$$

+ $N_2(v=0) + H_2O(001)$

(b)
$$3^{-13}$$

48)
$$N_2(v=1) + H_2O(000)$$

+ $N_2(v=0) + H_2O(010)$

(n)
$$9.4^{-11} \exp(-64/T^{1/3})$$

49)
$$H_2^{O(010)} + H_2^{O(010)}$$

$$+ \text{ H}_2\text{O(020)} + \text{H}_2\text{O(000)}$$
 (b) $1^{-12} \sqrt{\text{T}}$

50)
$$CO_2(001) + H_2O(000)$$

$$+ co_2(010) + H_2O(010)$$
 (b) 1^{-13}

51)
$$N_{2}(v=1) + CO_{2}(000)$$
 $+ N_{2}(v=0) + CO_{2}(001)$ (r) Take the largest value $\begin{cases} 1.71^{-6} \exp(-175.3/T^{1/3}) \\ 6.07^{-14} \exp(15.3/T^{1/3}) \end{cases}$

52) $N_{2}(v=1) + CO(v=0)$
 $+ N_{2}(v=0) - CO(v=1)$ (v) Take the largest value $\begin{cases} 1.78^{-6} \exp(-210/T^{1/3}) \\ 6.98^{-13} \exp(-25.6/T^{1/3}) \end{cases}$

53) $CO_{2}(001) + CO(v=0)$
 $+ CO_{2}(000) + CO(v=1)$ (b) $1.56^{-11} \exp(-30.1/T^{1/3})$

54) $CO_{2}(101) + H_{2}O(000)$
 $+ CO_{2}(000) + H_{2}O(001)$ (b) $2^{-13} \sqrt{T}$

55) $CO_{2}(021) + H_{2}O(000)$
 $+ CO_{2}(100) + CO_{2}(001)$ (b) $2^{-11} \sqrt{T}$

56) $CO_{2}(101) + CO_{2}(000)$
 $+ CO_{2}(100) + CO_{2}(010)$ (b) $4^{-13} \sqrt{T}$

57) $CO_{2}(020) + CO_{2}(000)$
 $+ CO_{2}(010) + CO_{2}(010)$ (b) $4^{-13} \sqrt{T}$

58) $H_{2}O(020) + H_{2} + H_{2}O(010) + H_{2}$ (w) 2^{-11}

60) $H_{2}O(020) + H_{2} + H_{2}O(010) + H_{2}$ (w) 2^{-11}

61) $H_{2}O(020) + H_{2}O + H_{2}O(010) + H_{2}O(010)$ (w) 2^{-11}

62) $CO_{2}(020) + H_{2} + CO_{2}(010) + H_{2}O(010)$ (w) 2^{-11}

63) $CO_{2}(020) + H_{2} + CO_{2}(010) + H_{2}O(010)$ (w) 2^{-11}

64) $CO_{2}(020) + H_{2}O + CO_{2}(010) + H_{2}O(010)$ (w) 2^{-11}

66)
$$CO_2(020) + CO_2 + CO_2(010) + CO_2$$
 (w) $9.2^{-10} \exp(-77/T^{1/3})$

67)
$$CO_2(030) + H \rightarrow CO_2(020) + H$$
 (w) $1.2^{-11} \exp(-7/T^{1/3})$

68)
$$CO_2(030) + O \rightarrow CO_2(020) + O$$
 (w) $6.9^{-9} \exp(-77/T^{1/3})$

69)
$$CO_2(030) + H_2 + CO_2(020) + H_2$$
 (w) $2.3^{-11} \exp(-7/T^{1/3})$

70)
$$CO_2(030) + H_2O + CO_2(020) + H_2O$$
 (w) 3^{-11}

71)
$$CO_2(030) + CO_2 + CO_2(020) + CO_2$$
 (w) 1.4⁻⁹ $exp(-77/T^{1/3})$

72) ${\rm H_2O}(100)$ and ${\rm H_2O}(001)$ assumed in quasi-equilibrium at the local translational temeprature.

73)
$$CO_2(100) + M \rightarrow CO_2(C20) + M$$
 (x) 1^{-10}

References (Vibrational Relaxation Model)

- (a) R.F. Hiedner and V.V. Casper, Chem. Phys. letters 15, 179(1972).
- (b) estimate
- (c) J.H. Kieffer and R.W. Lutz, J. Chem. Phys. 44, 658 (1966); 45, 3888(1966).
- (d) C. Joffrin, J. Ducuing and J.P. Coffinet, Opt, Commun. 2, 245(1970).
- (e) past measurements involving water presumably measure both VV and VT processes.
- (f) W.D. Breshears and P.F. Bird, J. Chem. Phys. $\underline{48}$, 4768 (1968).
- (g) R.J. McNeal, M.E. Whitson Jr. and G.R. Cook, preprint of work submitted to Chem. Phys. Letters (1972). Aerospace Report SPL 3303, July(1972).
- (h) D.J. Eckstrom, Stanford Research Institute, unpublished results.
- (i) D.R. White, J. Chem. Phys. 46, 2016(1967).
- (j) estimate based on Marriott's cross section values see E.R. Fisher, AFCRL Special Report, 72-0539, August (1972).
- (k) estimate based on all available data
- (1) estimate based on analogy to HF VT relaxation
- (m) estimate based on the modified Rapp model, J. Chem. Phys. 40, 573(1964) Fisher unpublished results.
- (n) estimate based on data of C.W. Von Rosenberg, K.N.C. Bray and N.H. Pratt, J. Chem. Phys. <u>56</u>, 3230(1972) and 13th Symp. (Intern.) on combustion p. 89, (1971).
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- (p) unpublished data from AVCO.

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- (r) R. Taylor and S. Bitterman, Rev. Mod. Phys. 41, 26 (1969).
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- (t) since direct VT relaxation of $CO_2(v_3)$ to ground state has not been unambiguously determined withoug competing VV processes, $CO_2(v_3)$ was assumed to relax like CO(v=1).
- (u) estimate based on analogy with reaction 35).
- (v) estimate based on AVCO suggestion.
- (w) scaled from the analogous rate constant for exciting the lowest bending mode.
- (x) the Fermi resonance in ${\rm CO}_2$ is assumed to couple the (100) and (020) modes with a gas kinetic efficiency independent of collision partner.
- (y) estimated based on the rate for CO+H, under (o) above.

(3) HYDROCARBON MODEL

LUMPED PARAMETER IR REACTION SCHEME

Reac	tion	Rate Constant, cc/sec	Photon Yield
1.	$A + O \rightarrow R + OH$	$1.3 \times 10^{-10} e^{-4840/RT}$	
2.	A + OH + R + H ₂ O	$1.2 \times 10^{-11} e^{-850/RT}$	1
3.	$R + O \rightarrow Ol + OH$	1 × 10 ⁻¹¹	6
4.	R + OH → Ol + H ₂ O	1 × 10 ⁻¹¹	7
5.	Ol + O → (80%) R + CHO	1×10^{-11}	
	→ (20%) C ₂ H ₅ + CHO		
6.	$C_{2}^{H_{6}} + OH + C_{2}^{H_{5}} + H_{2}^{O}$	$1.8 \times 10^{-11} e^{-2480/RT}$	1
7.	$C_2^{H_5} + O \rightarrow C_2^{H_4} + OH$	1 × 10 ⁻¹¹	5
8.	$C_{2}^{H_{5}} + OH + C_{2}^{H_{4}} + H_{2}^{O}$	1 × 10 ⁻¹¹	5
9.	$C_2H_4 + O \rightarrow CH_3 + CHO$	$5 \times 10^{-12} e^{-1040/RT}$	
10.	$C_2H_4 + OH \rightarrow CH_3 + CH_2O$	$1 \times 10^{-11} e^{+900/RT}$	
11.	$CH_4 + OH \rightarrow CH_3 + H_2O$	$4.0 \times 10^{-12} e^{-5000/RT}$	1
12.	$CH_3 + O_2 \rightarrow CH_2O + OH$	1.7×10^{-13}	1
13.	CHO + O → CO + OH	5 × 10 ⁻¹¹	8
14.	CH ₂ O + OH → CHO + H ₂ O	$4 \times 10^{-14} e^{-1000/RT}$	2
15.	CH ₂ O + O → CHO + OH	$1 \times 10^{-12} e^{-1000/RT}$	1
16.	сно + он →со + н ₂ о	1 × 10 ⁻¹²	9
17.	он + о → о ₂ + н	2.2×10^{-11}	

(3) HYDROCARBON MODEL

(cont'd)

18.	он + со + со ₂ + н	$9.3 \times 10^{-13} e^{-1080/RT}$	
19.	он + н ₂ + н ₂ о + н	$3.7 \times 10^{-11} e^{-5150/RT}$	1
20.	он + он → н ₂ о + о	$9.5 \times 10^{-12} e^{-780/RT}$	1
21.	он + н → н ₂ → о	$1.2 \times 10^{-11} e^{-7300/RT}$	
22.	н + 0 ₂ + он + о	3.7×10^{-10} e ^{-16,800/RT}	
23.	н + co ₂ + co + он	$9.3 \times 10^{-11} e^{-23,500/RT}$	
24.	н + н ₂ о + он + н ₂	1.4×10^{-9} e ^{-20,100/RT}	
25.	о + н ₂ о → он + он	9.5×10^{-11} e ^{-18,000/RT}	
26.	о + н ₂ + он + н	$2.9 \times 10^{-11} e^{-9450/RT}$	
27.	СH ₃ + 0 + CH ₂ 0 + H	3.3×10^{-11}	
	Сно + н + со + н ₂	3.3×10^{-11}	

"A" is a saturated hydrocarbon, $C_n^{H}_{2n+2}$, with $n \ge 3$

"R" is an alkyl radical, C_nH_{2n+1} , with $n \ge 3$

"Ol" is an olefinic hydrocarbon, $C_n^H_{2n}$, with $n \ge 3$

INFRARED HYDRAZINE MODEL A

FUEL CHEMI-EXCITATION REACTIONS

Reaction	Forward Rate Constant cc/sec	Reaction Energy kcal/gmole
$O + MMH \rightarrow HN_2CH_3 + H_2O$ (001) (100) H_2O (020) (010)	2 × 10 ⁻¹¹	-102
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1.8 × 10 ⁻¹¹	-90
O + UDMH \rightarrow HN ₂ (CH ₃) ₂ + OH OH + N ₂ H ₄ \rightarrow NH ₃ + NH ₂ O	1×10^{-12} 2×10^{-11}	-25

(cont'd)

REACTIONS OF OH AND H2O

Reaction

$$0 + OH \rightarrow 0_2 + H$$

$$OH + H \rightarrow H_2 + O$$

$$H_2^* + O \rightarrow OH + H$$

$$OH + H \rightarrow H_2^* + O$$

$$H_2^* + O \rightarrow OH^* + H$$

$$OH^* + H \rightarrow H_2^* + O$$

$$OH + H_2 \rightarrow H_2O$$

$$H_2O + H \rightarrow OH + H_2$$

$$OH + H_2^* \rightarrow H_2O^* + H$$

$$H_2O^* + H \to OH + H_2^*$$

$$OH^* + H \rightarrow OH + H_2$$

$$H_2O^* + H \rightarrow OH^* + H_2$$

$$OH + OH \rightarrow H_2O + O$$

$$O + OH + O_2 + H$$

$$0_2 + H \rightarrow OH + O$$

$$0 + OH^* \rightarrow O_2 + H$$

$$0_2 + H \rightarrow 0 + OH^*$$

Rate Constant cc/molecule-sec

$$4.0 \times 10^{-11}$$

$$2.9 \times 10^{-11} e^{-4725/T}$$

$$1.22 \times 10^{-11} e^{-3650/T}$$

$$0.65 \times 10^{-10}$$

$$2.75 \times 10^{-11} e^{-4925/T}$$

$$0.65 \times 10^{-10}$$

$$2.75 \times 10^{-11} e^{215/T}$$

$$3.8 \times 10^{-11} e^{-2600/T}$$

$$1.5 \times 10^{-10} e^{-10075/T}$$

$$1.5 \times 10^{-10}$$

$$5.8 \times 10^{-10} e^{-13475/T}$$

$$1.5 \times 10^{-11}$$

$$5.8 \times 10^{-11} e^{-12475/T}$$

$$9.6 \times 10^{-12} e^{-390/T}$$

$$9.5 \times 10^{-11} e^{-9000/T}$$

$$4 \times 10^{-11}$$

$$6.9 \times 10^{-10} e^{-8400/T}$$

$$4 \times 10^{-11}$$

$$6.9 \times 10^{-10} e^{-13400/T}$$

(4) AMINE MODEL (cont'd)

RELAXATION PROCESSES OF OH AND H2O

Reaction		Forward Rate Constant (cc/molecule-sec)
OH* + M ≠ OH + M	M = H	3×10^{-11}
II	M = O	3×10^{-11}
н	$M = H_2$	1×10^{-11}
H	$M = H_2O$	3×10^{-12}
$H_2O(010) + M \ddagger H_2O(000) + M$	M = H	2×10^{-10}
II	M = O	1×10^{-11}
ш	$M = H_2$	1×10^{-11}
ï	$M = H_2^0$	1 × 10 ⁻¹¹
$H_2^0(100) + M \rightarrow H_2^0 + M$	M = H	1 × 10 ⁻¹¹
u.	M = O	1×10^{-11}
U	$M = H_2$	2×10^{-13}
п	$M = H_2O$	1 × 10 ⁻¹²
$H_2O(001) + M + H_2O + M$	M = H	1 × 10 ⁻¹¹
11	M = O	1 × 10 ⁻¹¹
н	$M = H_2$	2 × 10 ⁻¹³
in.	$M = H_2O$	1 × 10 ⁻¹²

(4) AMINE MODEL (cont'd)

$H_2O(001) + M \rightarrow H_2O(020) + M$	M = H	1×10^{-11}
	M = O	1×10^{-11}
п	$M = H_2$	1×10^{-11}
an e	$M = H_2O$	1 × 10 ⁻¹¹
$H_2O(100) + M \rightarrow H_2O(020) + M$	M = H	1×10^{-11}
и	M = O	1×10^{-11}
a a	$M = H_2$	1×10^{-11}
u	$M = H_2O$	1×10^{-11}
$H_2O(010) + H_2O(010) + H_2O(020)$	+ H ₂ O	$1 \times 10^{-12} \sqrt{\mathrm{T}}$
$H_2O(020) + M \rightarrow H_2O(010) + M$	M = H	2×10^{-10}
**	M = O	1×10^{-11}
ıı .	$M = H_2$	1×10^{-11}
n i	$M = H_2O$	1×10^{-11}
$H + N_2H_4 + N_2H_3 + H_2$		$2.2 \times 10^{-11} e^{-1260/T}$
$H + N_2H_3 \rightarrow 2NH_2$		
$H + N_2H_2 \rightarrow N_2H_3$		1 × 10 ⁻¹⁰
$O + NH_3 + NH_2 + OH$		$2.5 \times 10^{-12} e^{-3000/T}$
$OH + NH_3 + NH_2 + H_2O^{(*)}$		$6.6 \times 10^{-14} e^{-560/T}(T).68$

(5) RADIATIVE LIFETIMES

SPECIES	t(sec)
OH (v=1)	0.3
H ₂ 0(100)	0.14
H ₂ 0(010)	0.045
H ₂ 0(001)	0.014
H ₂ 0(020)	0.022
CO ₂ (010)	0.33
CO ₂ (001)	0.0025
CO ₂ (101)	0.091